

# **BOILING**

---

**Nazaruddin Sinaga**

**Efficiency and Energy Conservation Laboratory  
Diponegoro University**

# General Considerations

- Boiling is associated with **transformation of liquid to vapor at a solid/liquid interface** due to convection heat transfer from the solid.
- Agitation of fluid by vapor bubbles provides for **large convection coefficients** and hence **large heat fluxes** at **low-to-moderate surface-to-fluid temperature differences**.
- Special form of **Newton's law of cooling**:

$$q_s'' = h(T_s - T_{sat}) = h \Delta T_e$$

➤  $T_{sat}$  → **saturation temperature** of liquid

➤  $\Delta T_e \equiv (T_s - T_{sat})$  → **excess temperature**

- Special Cases

- **Pool Boiling:**

- Liquid motion is due to natural convection and bubble-induced mixing.

- **Forced Convection Boiling:**

- Fluid motion is induced by external means, as well as by bubble-induced mixing.

- **Saturated Boiling:**

- Liquid temperature is slightly larger than saturation temperature.

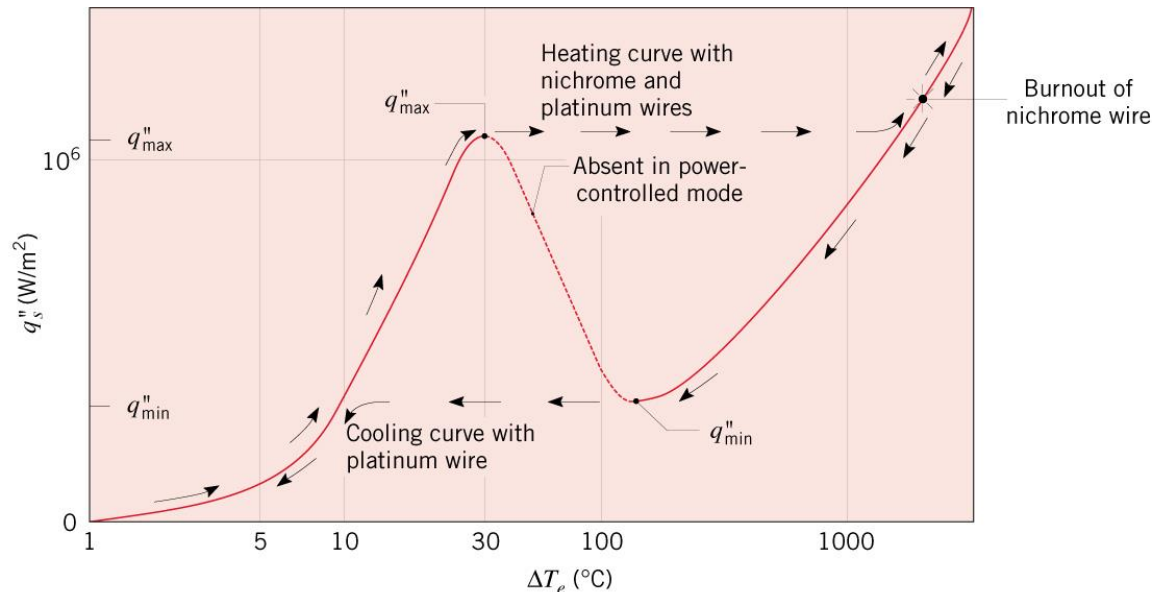
- **Subcooled Boiling:**

- Liquid temperature is less than saturation temperature.

# The Boiling Curve

Reveals range of conditions associated with **saturated pool boiling** on a  $q_s'' - \Delta T_e$  plot.

Water at Atmospheric Pressure

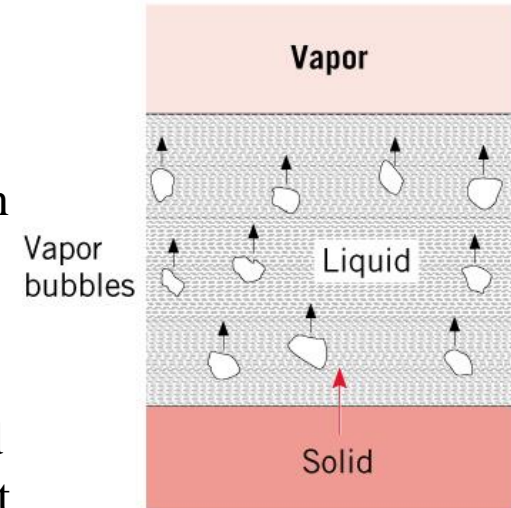


- **Free Convection Boiling** ( $\Delta T_e < 5^\circ C$ )
  - Little vapor formation.
  - Liquid motion is due principally to single-phase natural convection.
- **Onset of Nucleate Boiling - ONB** ( $\Delta T_e \approx 5^\circ C$ )

- **Nucleate Boiling ( $5 < \Delta T_e < 30^\circ C$ )**

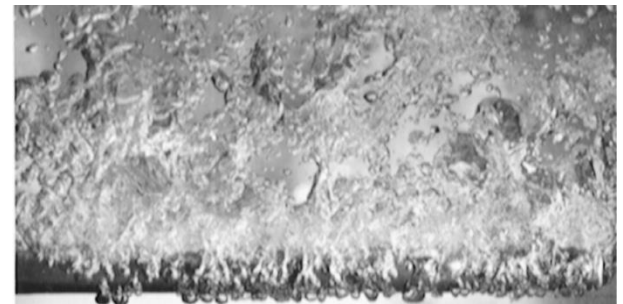
- **Isolated Vapor Bubbles ( $5 < \Delta T_e < 10^\circ C$ )**

- Liquid motion is strongly influenced by nucleation of bubbles at the surface.
- $h$  and  $q_s''$  increase sharply with increasing  $\Delta T_e$ .
- Heat transfer is principally due to contact of liquid with the surface (single-phase convection) and not to vaporization.

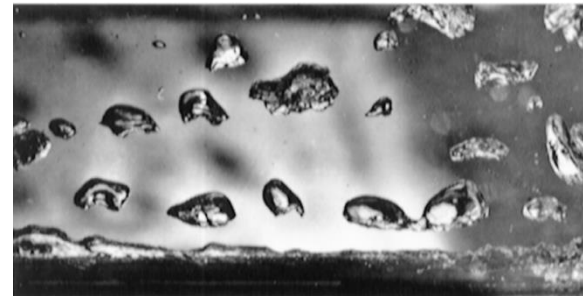


- **Jets and Columns ( $10 < \Delta T_e < 30^\circ C$ )**

- Increasing number of nucleation sites causes bubble interactions and coalescence into jets and slugs.
- Liquid/surface contact is impaired.
- $q_s''$  continues to increase with  $\Delta T_e$  while  $h$  begins to decrease.



- **Critical Heat Flux - CHF,  $q''_{\max}$  ( $\Delta T_e \approx 30^\circ C$ )**
  - Maximum attainable heat flux in nucleate boiling.
  - $q''_{\max} \approx 1 \text{ MW/m}^2$  for water at atmospheric pressure.
- **Potential Burnout for Power-Controlled Heating**
  - An increase in  $q''_s$  beyond  $q''_{\max}$  causes the surface to be blanketed by vapor, and the surface temperature can spontaneously achieve a value that potentially exceeds its melting point ( $\Delta T_s > 1000^\circ C$ ).
  - If the surface survives the temperature shock, conditions are characterized by **film boiling**.
- **Film Boiling**
  - Heat transfer is by conduction and radiation across the **vapor blanket**.
  - A reduction in  $q''_s$  follows the cooling curve continuously to the **Leidenfrost point** corresponding to the **minimum heat flux  $q''_{\min}$  for film boiling**.



- A reduction in  $q_s''$  below  $q_{\min}''$  causes an abrupt reduction in surface temperature to the nucleate boiling regime.
- **Transition Boiling for Temperature-Controlled Heating**
  - Characterized by a continuous decay of  $q_s''$  (from  $q_{\max}''$  to  $q_{\min}''$ ) with increasing  $\Delta T_e$ .
  - Surface conditions oscillate between nucleate and film boiling, but portion of surface experiencing film boiling increases with  $\Delta T_e$ .
  - Also termed **unstable** or **partial film boiling**.

# Pool Boiling Correlations

- Nucleate Boiling

- Rohsenow Correlation

$$q_s'' = \mu_l h_{fg} \left[ \frac{g(\rho_l - \rho_v)}{\sigma} \right]^{1/2} \left( \frac{c_{p,l} \Delta T_e}{C_{s,f} h_{fg} \text{Pr}_l^n} \right)^3$$

$C_{s,f}, n \rightarrow$  Surface/Fluid Combination (Table 10.1)

- Critical Heat Flux

$$q_{\max}'' = 0.149 h_{fg} \rho_v \left[ \frac{\sigma g (\rho_l - \rho_v)}{\rho_v^2} \right]^{1/4}$$



- Film Boiling**

The cumulative (and coupled effects) of convection and radiation across the vapor layer  $\longrightarrow$

$$\bar{h}^{4/3} \approx \bar{h}_{conv}^{4/3} + \bar{h}_{rad} \bar{h}^{1/3}$$

$$\overline{Nu}_D = \frac{\bar{h}_{conv} D}{k_v} = C \left[ \frac{g(\rho_l - \rho_v) h'_{fg} D^3}{\nu_v k_v (T_s - T_{sat})} \right]^{1/4}$$

Geometry	C
Cylinder(Hor.)	0.62
Sphere	0.67

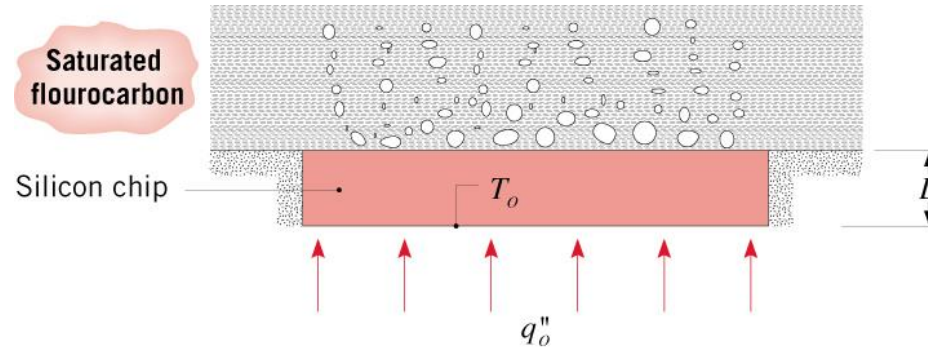
$$h'_{fg} = h_{fg} + 0.80 c_{p,v} (T_s - T_{sat})$$

$$\bar{h}_{rad} = \frac{\varepsilon \sigma (T_s^4 - T_{sat}^4)}{T_s - T_{sat}}$$

If  $\bar{h}_{conv} > \bar{h}_{rad}$ ,

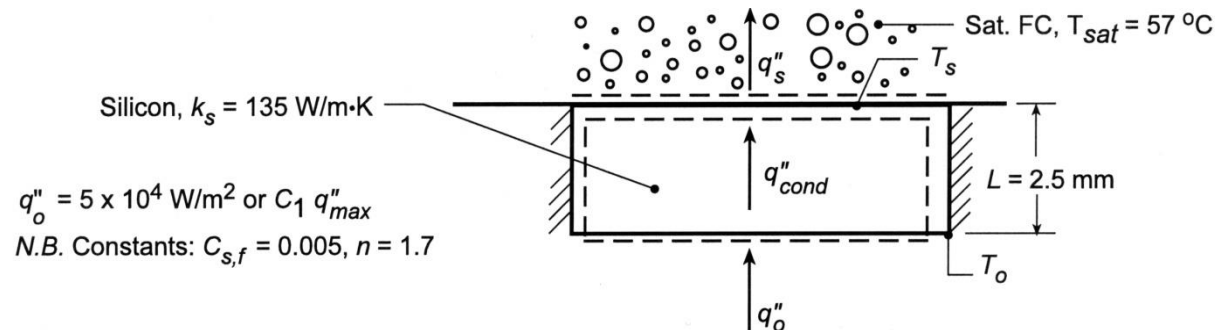
$$\bar{h} \approx \bar{h}_{conv} + 0.75 \bar{h}_{rad}$$

## Problem 10.23: Chip thermal conditions associated with cooling by immersion in a fluorocarbon.



**KNOWN:** Thickness and thermal conductivity of a silicon chip. Properties of saturated fluorocarbon liquid

**FIND:** (a) Temperature at bottom surface of chip for a prescribed heat flux and for a flux that is 90% of CHF, (b) Effect of heat flux on chip surface temperatures; maximum allowable heat flux for a surface temperature of  $80^\circ\text{C}$ .



**ASSUMPTIONS:** (1) Steady-state conditions, (2) Uniform heat flux and adiabatic sides, hence one-dimensional conduction in chip, (3) Constant properties, (4) Nucleate boiling in liquid.

**PROPERTIES:** Saturated fluorocarbon (given):  $c_{p,\ell} = 1100 \text{ J/kg}\cdot\text{K}$ ,  $h_{fg} = 84,400 \text{ J/kg}$ ,  $\rho_\ell = 1619.2 \text{ kg/m}^3$ ,  $\rho_v = 13.4 \text{ kg/m}^3$ ,  $\sigma = 8.1 \times 10^{-3} \text{ kg/s}^2$ ,  $\mu_\ell = 440 \times 10^{-6} \text{ kg/m}\cdot\text{s}$ ,  $\text{Pr}_\ell = 9.01$ .

**ANALYSIS:** (a) Energy balances at the top and bottom surfaces yield  $q_o'' = q_{\text{cond}}'' = k_s (T_o - T_s)/L = q_s''$ ; where  $T_s$  and  $q_s''$  are related by the Rohsenow correlation,

$$T_s - T_{\text{sat}} = \frac{C_{s,f} h_{fg} \text{Pr}_\ell^n}{c_{p,\ell}} \left( \frac{q_s''}{\mu_\ell h_{fg}} \right)^{1/3} \left[ \frac{\sigma}{g(\rho_\ell - \rho_v)} \right]^{-1/6}$$

Hence, for  $q_s'' = 5 \times 10^4 \text{ W/m}^2$ ,

$$T_s - T_{\text{sat}} = \frac{0.005(84,400 \text{ J/kg})9.01^{1.7}}{1100 \text{ J/kg}\cdot\text{K}} \left( \frac{5 \times 10^4 \text{ W/m}^2}{440 \times 10^{-6} \text{ kg/m}\cdot\text{s} \times 84,400 \text{ J/kg}} \right)^{1/3} \\ \times \left[ \frac{8.1 \times 10^{-3} \text{ kg/s}^2}{9.807 \text{ m/s}^2 (1619.2 - 13.4) \text{ kg/m}^3} \right]^{-1/6} = 15.9^\circ \text{C}$$

$$T_s = (15.9 + 57)^\circ \text{C} = 72.9^\circ \text{C}$$

From the rate equation,

$$T_o = T_s + \frac{q_o'' L}{k_s} = 72.9^\circ\text{C} + \frac{5 \times 10^4 \text{ W/m}^2 \times 0.0025 \text{ m}}{135 \text{ W/m}\cdot\text{K}} = 73.8^\circ\text{C}$$

For a heat flux which is 90% of the critical heat flux ( $C_1 = 0.9$ ),

$$q_o'' = 0.9q_{\text{max}}'' = 0.9 \times 0.149 h_{\text{fg}} \rho_v \left[ \frac{\sigma g (\rho_l - \rho_v)}{\rho_v^2} \right]^{1/4} = 0.9 \times 0.149 \times 84,400 \text{ J/kg} \times 13.4 \text{ kg/m}^3$$

$$\times \left[ \frac{8.1 \times 10^{-3} \text{ kg/s}^2 \times 9.807 \text{ m/s}^2 (1619.2 - 13.4) \text{ kg/m}^3}{(13.4 \text{ kg/m}^3)^2} \right]^{1/4}$$

$$q_o'' = 0.9 \times 15.5 \times 10^4 \text{ W/m}^2 = 13.9 \times 10^4 \text{ W/m}^2$$

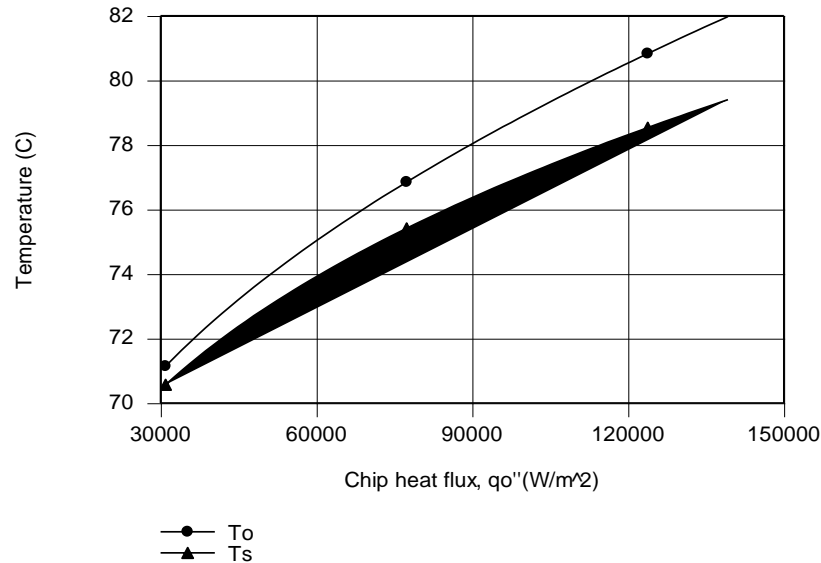
From the results of the previous calculation and the Rohsenow correlation, it follows that

$$\Delta T_e = 15.9^\circ\text{C} \left( q_o'' / 5 \times 10^4 \text{ W/m}^2 \right)^{1/3} = 15.9^\circ\text{C} (13.9/5)^{1/3} = 22.4^\circ\text{C}$$

Hence,  $T_s = 79.4^\circ\text{C}$  and

$$T_o = 79.4^\circ\text{C} + \frac{13.9 \times 10^4 \text{ W/m}^2 \times 0.0025 \text{ m}}{135 \text{ W/m}\cdot\text{K}} = 82^\circ\text{C}$$

(b) Parametric calculations for  $0.2 \leq C_1 \leq 0.9$  yield the following variations of  $T_s$  and  $T_o$  with  $q_o''$ .



The chip surface temperatures, as well as the difference between temperatures, increase with increasing heat flux. The maximum chip temperature is associated with the bottom surface, and  $T_o = 80^\circ\text{C}$  corresponds to

$$q_{o,\max}'' = 11.3 \times 10^4 \text{ W/m}^2 \quad <$$

which is 73% of CHF ( $q_{\max}'' = 15.5 \times 10^4 \text{ W/m}^2$ ).

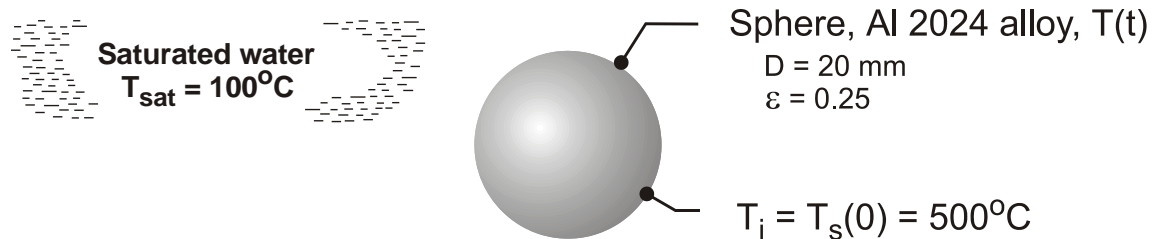
**COMMENTS:** Many of today's VLSI chip designs involve heat fluxes well in excess of  $15 \text{ W/cm}^2$ , in which case pool boiling in a fluorocarbon would not be an appropriate means of heat dissipation.

**Problem 10.26:** Initial heat transfer coefficient for immersion of an aluminum sphere in a saturated water bath at atmospheric pressure and its temperature after immersion for 30 seconds.

**KNOWN:** A sphere (aluminum alloy 2024) with a uniform temperature of  $500^{\circ}\text{C}$  and emissivity of 0.25 is suddenly immersed in a saturated water bath maintained at atmospheric pressure.

**FIND:** (a) The total heat transfer coefficient for the initial condition; fraction of the total coefficient contributed by radiation; and (b) Estimate the temperature of the sphere 30 s after it has been immersed in the bath.

### SCHEMATIC



**ASSUMPTIONS:** (1) Water is at atmospheric pressure and uniform temperature,  $T_{\text{sat}}$ , and  
(2) Lumped capacitance method is valid.

**PROPERTIES:**

Vapor Blanket ( $T_{f,i} = 573\text{K}$ ):  $h_{fg} = 1.41 \times 10^6 \text{ J/kg}$ ,  $k_v = 0.0767 \text{ W/m} \cdot \text{K}$ ,  $c_{p,v} = 5889 \text{ J/kg} \cdot \text{K}$ ,  
 $\text{Pr}_v = 1.617$ ,  $\rho_v = 46.0 \text{ kg/m}^3$ ,  $\nu_v = 4.33 \times 10^{-7} \text{ m}^2/\text{s}$ ,  $\rho_l = 712 \text{ kg/m}^3$ .

Aluminum Alloy:  $\rho_s = 2700 \text{ kg/m}^3$ ,  $c_{p,s} = 875 \text{ J/kg} \cdot \text{K}$ ,  $k_s = 186 \text{ W/m} \cdot \text{K}$ .

**ANALYSIS:** (a) For the initial condition with  $T_s = 500^\circ\text{C}$ , *film boiling* will occur and the coefficients due to convection and radiation are estimated using Eqs. 10.9 and 10.11, respectively,

$$\overline{\text{Nu}}_D = \frac{\bar{h}_{\text{conv}} D}{k_v} = C \left[ \frac{g(\rho_l - \rho_v) h'_{fg} D^3}{\nu_v k_v (T_s - T_{\text{sat}})} \right]^{1/4} \quad (1)$$

$$\bar{h}_{\text{rad}} = \frac{\varepsilon \sigma (T_s^4 - T_{\text{sat}}^4)}{T_s - T_{\text{sat}}} \quad (2)$$

where  $C = 0.67$  for spheres and  $\sigma = 5.67 \times 10^{-8} \text{ W/m}^2 \cdot \text{K}^4$ . The corrected latent heat is

$$h'_{fg} = h_{fg} + 0.8 c_{p,v} (T_s - T_{\text{sat}}) \quad (3)$$

The total heat transfer coefficient is given by Eq. 10.10a as

$$\bar{h}^{4/3} = \bar{h}_{\text{conv}}^{4/3} + \bar{h}_{\text{rad}} \cdot \bar{h}^{1/3} \quad (4)$$

Using the foregoing relations, the following results are obtained.

$$\begin{array}{cccc} \bar{\text{Nu}}_D & \bar{h}_{\text{cnv}} \left( \text{W/m}^2 \cdot \text{K} \right) & \bar{h}_{\text{rad}} \left( \text{W/m}^2 \cdot \text{K} \right) & \bar{h} \left( \text{W/m}^2 \cdot \text{K} \right) \\ 226 & 867 & 12.0 & 876 \end{array}$$

(b) For the lumped-capacitance method, from Section 5.3, the energy balance is

$$-\bar{h}A_s (T_s - T_{\text{sat}}) = \rho_s V c_s \frac{dT_s}{dt} \quad (5)$$

where  $\rho_s$  and  $c_s$  are properties of the sphere. Numerically integrating Eq. (5) and evaluating  $\bar{h}$  as a function of  $T_s$ , the following result is obtained for the sphere temperature after 30s.

$$T_s(30\text{s}) = 333^\circ\text{C}.$$

**COMMENTS:** (1) The Biot number associated with the aluminum alloy sphere cooling process for the initial condition is  $\text{Bi} = 0.09$ . Hence, the lumped-capacitance method is valid.

(2) Radiation makes a negligible contribution to the heat rate throughout the process.